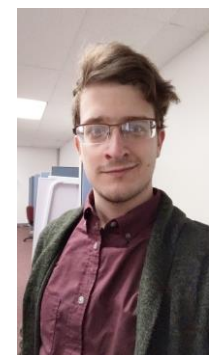


Large Magnetoresistance and Quantum Oscillations in $\text{Sn}_{0.05}\text{Pb}_{0.95}\text{Te}$



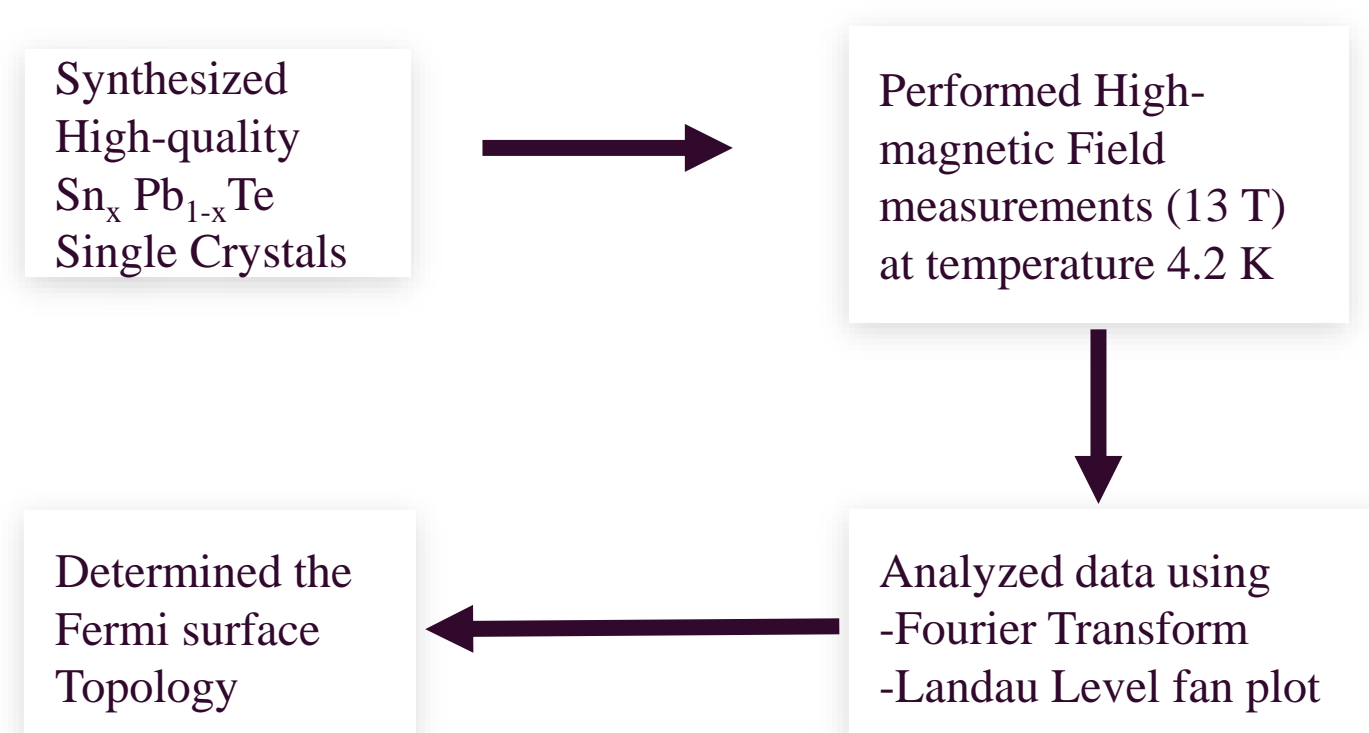
PRESENTER

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Background

Topological Insulators are very new and interesting materials due to their highly conducting surface states arising from the non-trivial bulk topology. They have the potential to be used in a variety of areas such as highly efficient electronics and possibly quantum computers. Here, we have studied magnetotransport properties of $\text{Sn}_x\text{Pb}_{1-x}\text{Te}$, known as a Topological Crystalline Insulator.

Methods



Results

Figure 2 describes longitudinal (R_{xx}) and Hall resistance (R_{xy}) under high magnetic fields up to 13 T at 4.2 K. At high fields, both R_{xx} and R_{xy} show oscillations, known as Shubnikov de-Haas (SdH) oscillations. The oscillations are clearly visible in the background subtracted data (insets).

Figure 3 shows the frequencies of SdH oscillations at different tilt angle (θ) of the magnetic field. There exists mainly two frequencies (f_α and f_β).

Figure 4 shows the Landau level fan plot for $\text{Sn}_{0.05}\text{Pb}_{0.95}\text{Te}$ single crystal. The maxima and minima positions are assigned as integer and half-integer numbers, respectively. From the linear extrapolation, we observed the Berry phase ~ 0.1 , which is close to 0 (expected for a topological trivial system).

- A topological Insulators is a new quantum state of materials with an insulating bulk and highly conducting surface.
- The bulk has a trivial insulating bands, whereas the surface has a linear band structure as shown in Fig. 1.
- We synthesized high-quality single crystals of $\text{Sn}_x\text{Pb}_{1-x}\text{Te}$.
- Our magnetotransport data reveal the trivial topology (not a TI) of $\text{Sn}_{0.05}\text{Pb}_{0.95}\text{Te}$.

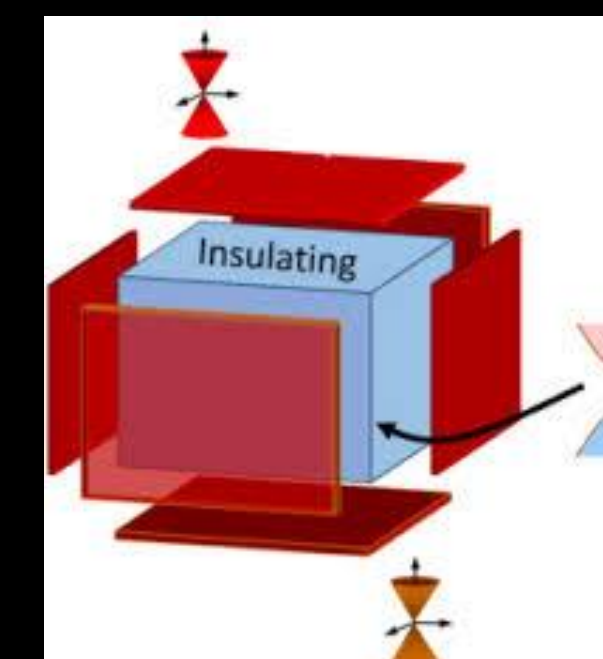


Fig. 1. A schematic diagram for a topological insulator showing the bulk and surface along with their band structure.

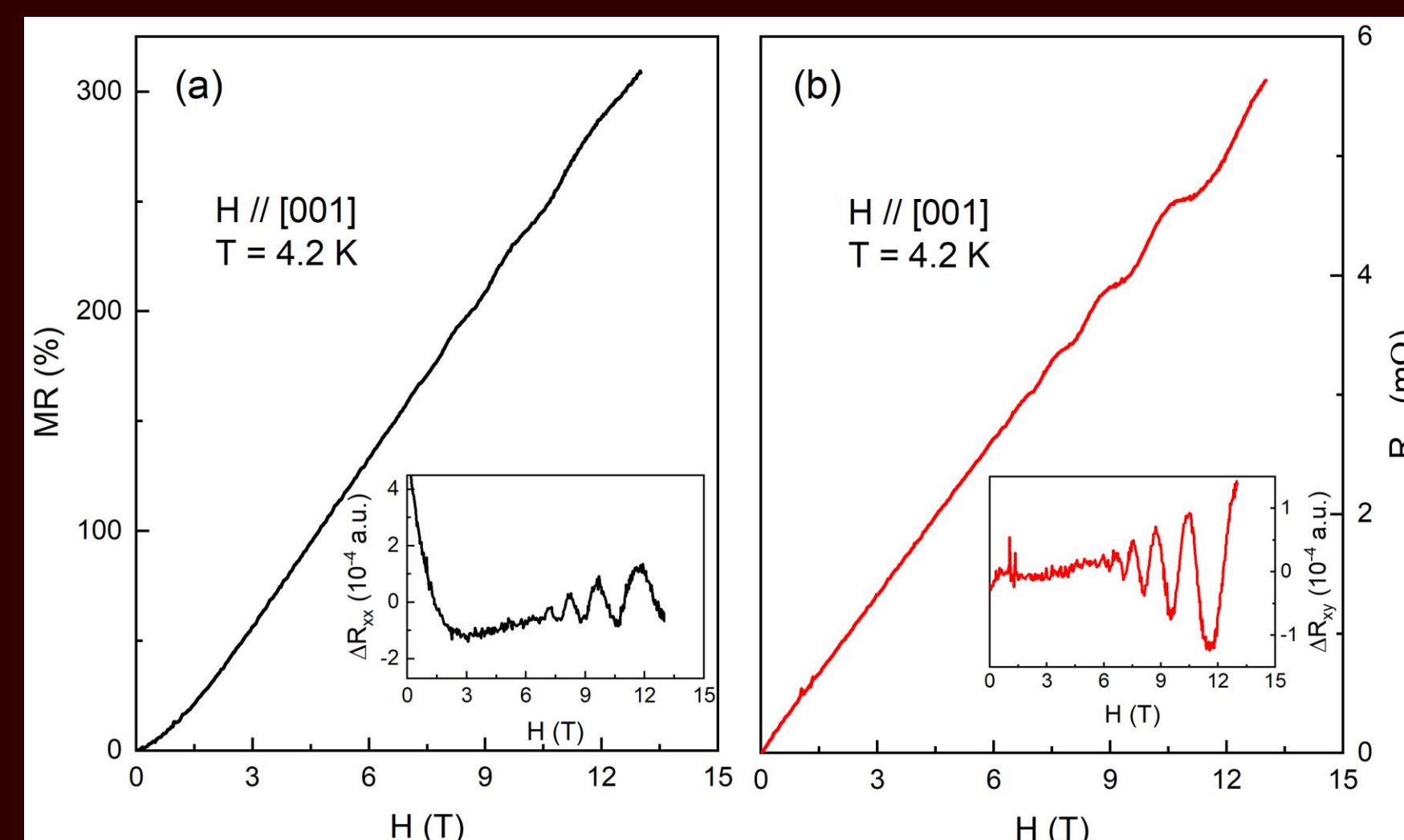


Figure 2: Magnetoresistance and Quantum Oscillations. (a) Magnetoresistance (MR) expressed as a percentage and (b) Hall resistance (R_{xy}) of single-crystal $\text{Sn}_{0.05}\text{Pb}_{0.95}\text{Te}$ as function of applied field (H). MR is positive and reaches $\sim 310\%$ under an applied field of 13 T at 4.2 K. The positive slope of the R_{xy} vs. H plot indicates the presence of hole-like bulk charge carriers. Both MR and R_{xy} show a signature of quantum oscillations at higher fields that is clearly observed in the polynomial background-subtracted data, as shown in the insets to (a) and (b), respectively. The oscillations are periodic and appear to have a single frequency.

Figure 3: Frequency Plot at Higher Angles. Fourier transforms of quantum oscillations for both longitudinal and Hall resistance at θ values of (a) 0° , (b) 15° , (c) 30° , and (d) 45° . The frequency values obtained from oscillations in longitudinal and transverse resistance are comparable to one another. At $\theta = 0^\circ$, there exists only one frequency. The number of frequencies increases at higher angles, suggesting the existence of more than one Fermi wave vector.

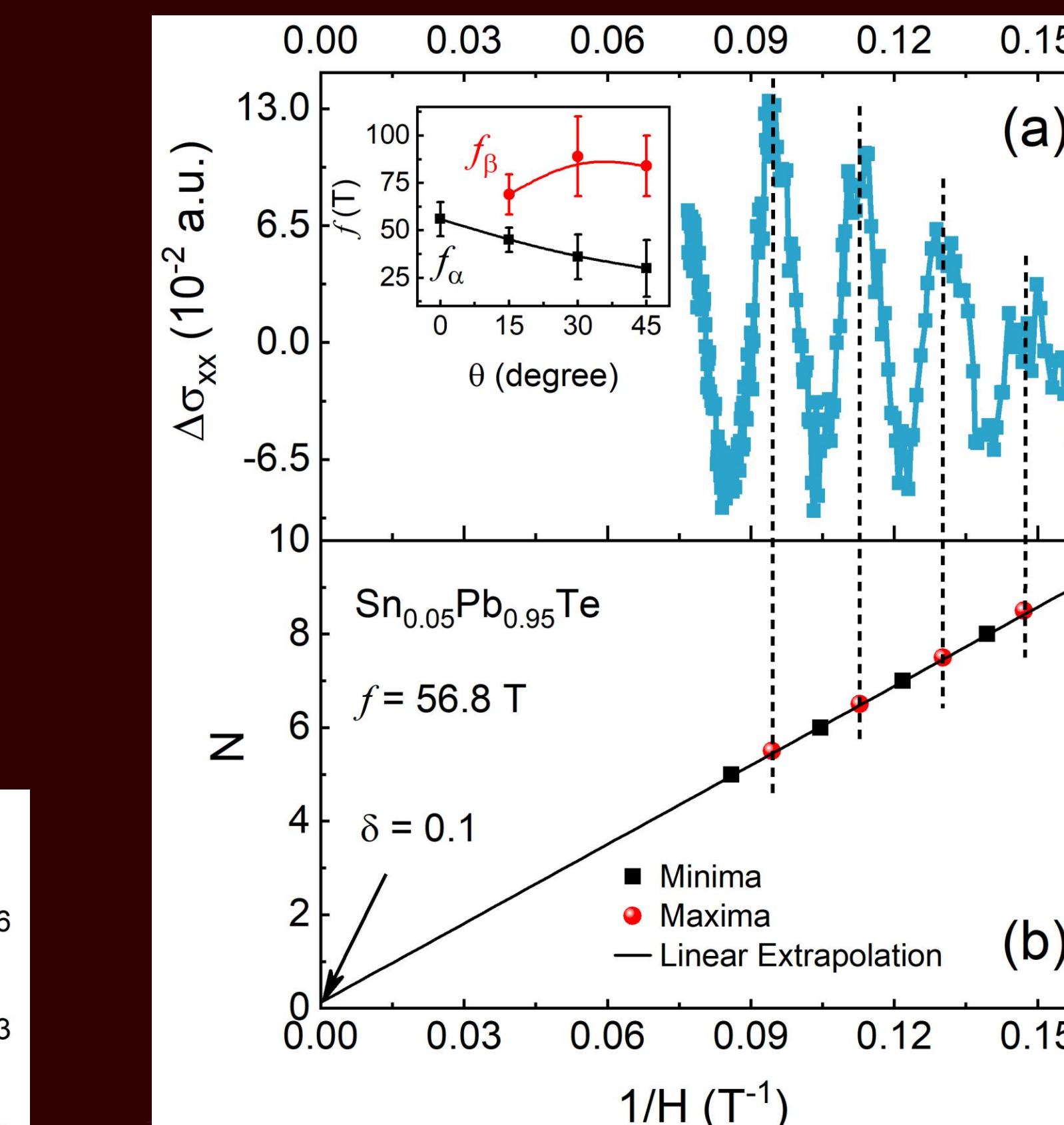
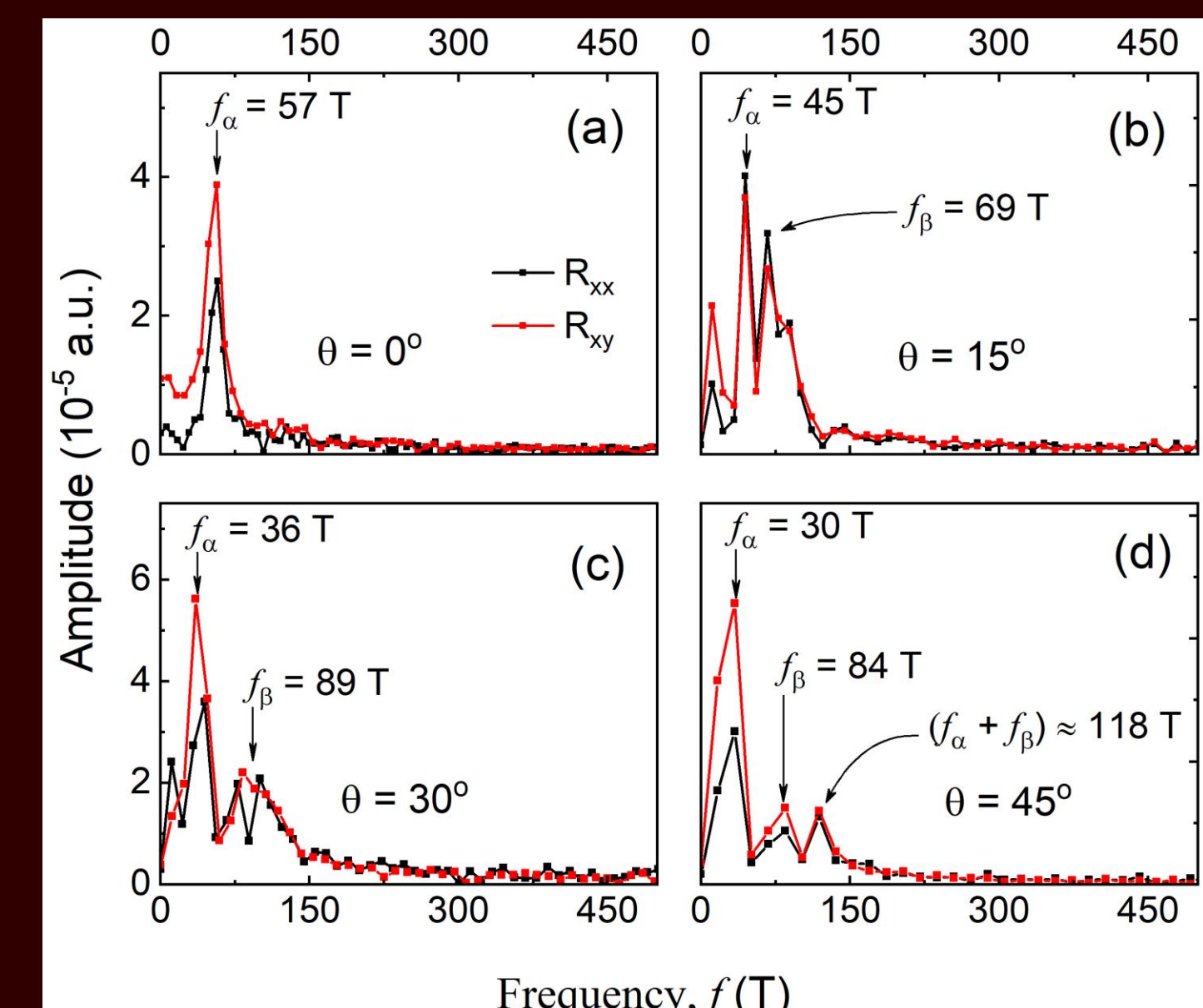


Figure 4: Berry Phase Analyses. (a) Quantum oscillations after background subtraction at $\theta = 0^\circ$. Inset: angle dependence of frequencies f_α and f_β from Fig. 5. The error bar for each data point is defined as the half-width at half height of the respective peak in the frequency plot [Fig. (3)]. (b) Landau-level fan diagram constructed using maxima and minima of the oscillations in (a). From the linear extrapolation, the Berry phase and frequency values are found to be $\delta = 0.1$ and $f_\alpha = 56.8$ T, respectively.

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